

Entropic Lattice Boltzmann Method: an implicit Large-Eddy Simulation?

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November 20, 2017



This project has received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement No' 642069 and was conducted within the activity of ERC Grant No' 339032

Motivations

Context

Implicit SGS arising when stabilizing for $\nu \rightarrow 0$ a Lattice Boltzmann Method by equipping it with a H-theorem

Aim

Study the physical properties of this implicit sub-grid scale (SGS) model

Tool

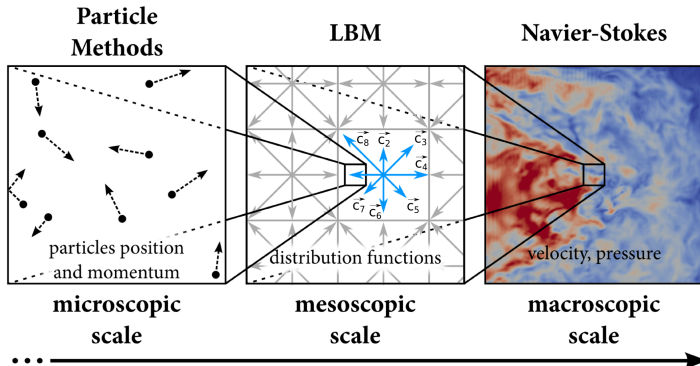
Development a tool based on balance equations to check hydrodynamic recovery of simulated flows accross scales

Introduction to Lattice Boltzmann Method (LBM)

LBM Equation with a relaxation time $\tau \equiv \tau_0$ fixed

$$f_i(\vec{x} + \vec{c}_i \Delta t, t + \Delta t) - f_i(\vec{x}, t) = -\frac{1}{\tau_0} [f_i(\vec{x}, t) - f_i^{eq}(\vec{x}, t)]$$

Macroscopic quantities: Density: $\rho = \sum_i f_i$ Momentum: $\rho \vec{u} = \sum_i f_i \vec{c}_i$



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Chapman-Enskog expansion: $\mathcal{O}(K_n^2)$, $\mathcal{O}(M_a^3)$

$$\nu = c_s^2 (\tau - 0.5) \Delta t$$



Weakly compressible Navier-Stokes with viscosity $\nu \equiv \nu_0$ fixed

$$\rho \partial_t \mathbf{u}_i + \rho u_j \partial_j \mathbf{u}_i = -\partial_i p + \partial_j \rho \nu (\partial_j \mathbf{u}_i + \partial_i \mathbf{u}_j) + \mathcal{O}(M_a^3) + \mathcal{O}(K_n^2)$$

with c_s^2 the speed of sound is the lattice

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Instabilities arising as $\tau_0 \rightarrow 0.5 \iff \nu_0 \rightarrow 0$ at a fixed resolution
Can we get rid of those instabilities?

Entropic LBM: Unconditionnal stability and implicit SGS

- ▶ Entropic LBM (ELBM) equips a H-theorem by locally adaptating τ

$$\tau = \tau^{eff}(\vec{x}, t) = \frac{2\tau_0}{\alpha(f_i(\vec{x}, t))} \text{ where } \alpha \text{ has a non-linear dependency on } f_i(\vec{x}, t)$$

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$$\alpha \approx 2 \iff \tau^{\text{eff}} \approx \tau_0$$

[Malaspinas *et al.*, 2008]

$$\nu_{\text{eff}}(\vec{x}, t) = \nu_0 + \nu_t^M(\vec{x}, t) \text{ with } \nu_t^M(\vec{x}, t) \propto -\frac{S_{\theta\kappa} S_{\kappa\gamma} S_{\gamma\theta}}{S_{\lambda\mu} S_{\lambda\mu}}$$

where $S_{ij} = \frac{1}{2}(\partial_i u_j + \partial_j u_i)$

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Similar to a Smagorinsky model: $\nu_t(\vec{x}, t) = C\sqrt{S_{\theta\kappa} S_{\theta\kappa}}$

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Objective: Numerically check the existence of this implied SGS

Kinetic energy balancing averaged over a sub-volume V

Averaged kinetic energy balance equation for $\nu = \nu_0$ fixed

$$\begin{aligned}LHS_V^E &= \partial_t \left\langle \frac{\rho u_i u_i}{2} \right\rangle_V \\&= - \left\langle u_i \partial_i p \right\rangle_V - \nu_0 \left\langle \rho (\partial_j u_i + \partial_i u_j) \partial_j u_i \right\rangle_V + \nu_0 \left\langle \partial_j \rho u_i (\partial_j u_i + \partial_i u_j) \right\rangle_V \\&\quad - \left\langle \partial_j \frac{\rho u_i u_i}{2} u_j \right\rangle_V + \left\langle u_i F_i \right\rangle_V \\&= RHS_V^{E,1} + RHS_V^{E,2} + RHS_V^{E,3} + RHS_V^{E,4} + RHS_V^{E,5} = RHS_V^E\end{aligned}$$

where $\langle \dots \rangle_V$ denotes the average over a sub-volume V

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Averaged kinetic energy balance equation for $\nu = \nu^{eff}(\vec{x}, t) = \nu_0 + \nu_1(\vec{x}, t)$

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 &\quad - \left\langle \partial_j \frac{\rho u_i u_i}{2} u_j \right\rangle_V + \langle u_i F_i \rangle_V - \langle \nu_1 \rho (\partial_j u_i + \partial_i u_j) \partial_j u_i \rangle_V + \langle \partial_j \nu_1 \rho u_i (\partial_j u_i + \partial_i u_j) \rangle_V \\
 &= RHS_V^{E,1} + RHS_V^{E,2} + RHS_V^{E,3} + RHS_V^{E,4} + RHS_V^{E,5} + RHS_V^{E,6} + RHS_V^{E,7} \\
 &= RHS_V^E
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Simulations of 2D homogeneous isotropic turbulence

Periodic 256×256 grid using a D2Q9 lattice

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Forcing on a shell of wavenumber

$$F_{\Psi}^T = F_0^T \sum_{\|\vec{k}\|=5}^7 \cos\left(\frac{2\pi}{L} \vec{k} \cdot \vec{x} + \phi\right)$$

where ϕ is an arbitrary constant

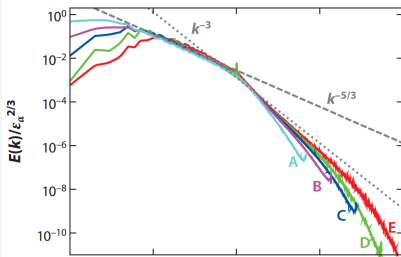
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When forcing at k_f , we have:

- ▶ a backward energy cascade to large scales
- ▶ a forward enstrophy cascade to small scales

[Boffetta & Ecke, 2012]

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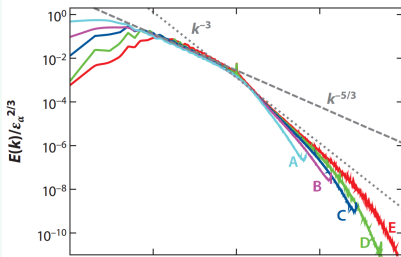
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Energy removal at large scale

$$\vec{F}^E(\vec{x}, t) = -F_0^E \sum_{\|\vec{k}\|=1}^2 \vec{u}(\vec{k}, t) e^{\frac{2\pi}{L} \vec{k} \cdot \vec{x}}$$



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Statistical analysis of the balancing errors δ_V^E

Kinetic energy balancing error

$$\delta_V^E(t) = \frac{|RHS_V^E(t) - LHS_V^E(t)|}{\max_i |RHS_V^{E,i}(t)|}$$

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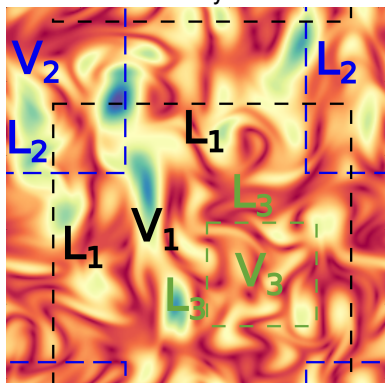
For a scale L , we gather statistics of δ_L^E :
Balancing error over sub-volumes of shape $V = L \times L$ in space and time

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Example of sub-volumes shown
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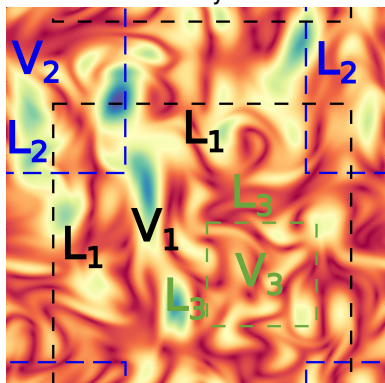
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We process outputs from simulations that have reached statistical stationarity

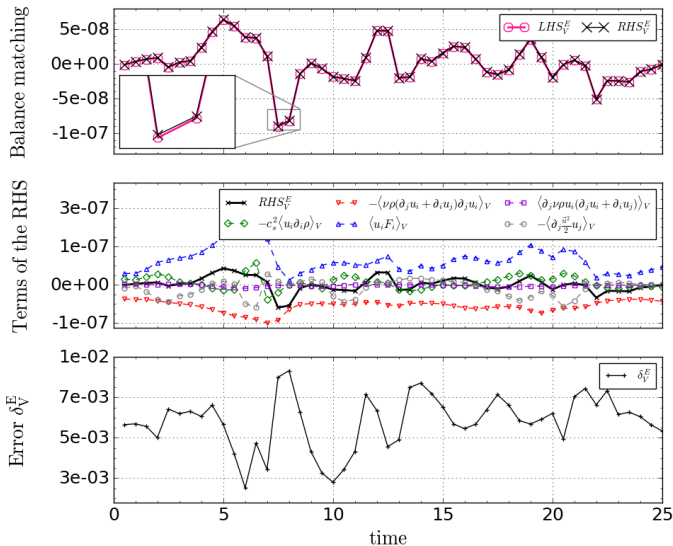
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Evolution of the kinetic energy balancing on a single sub-volume

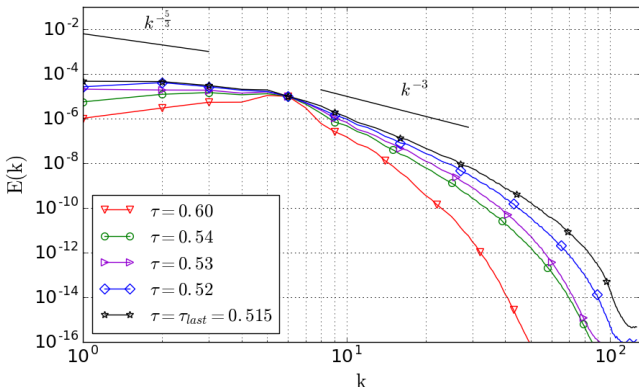
$V = 181 \times 181$ for a LBM simulation with $\tau \equiv \tau_0 = 0.55$ fixed



LBM for different fixed $\tau \equiv \tau_0 \rightarrow 0.5$: Superposed spectrum

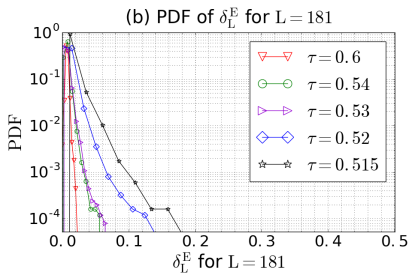
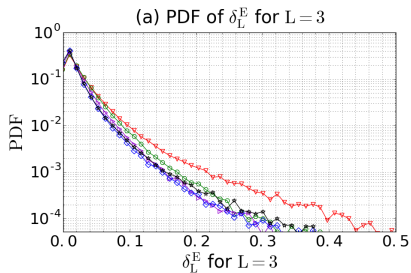
$\tau \equiv \tau_0$ fixed & Constant forcing amplitudes

Time-averaged superposed spectrum

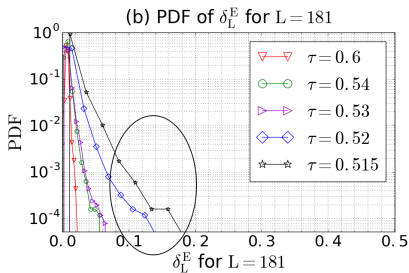
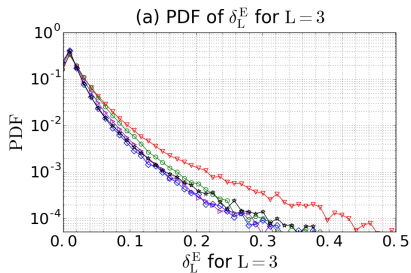


where $\tau_{last} = 0.515$ is the lowest τ ensuring a stable LBM simulation at this resolution and forcing amplitude

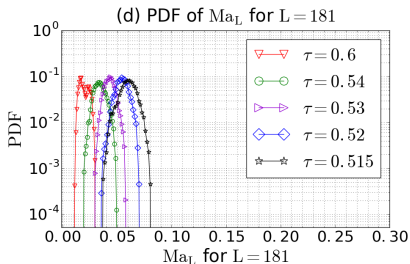
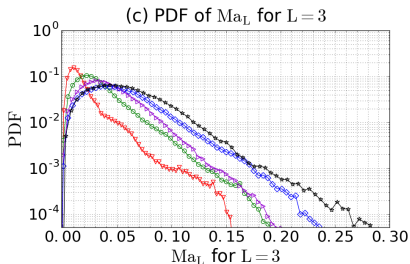
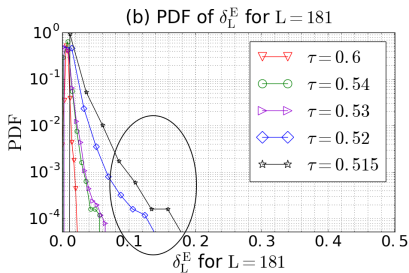
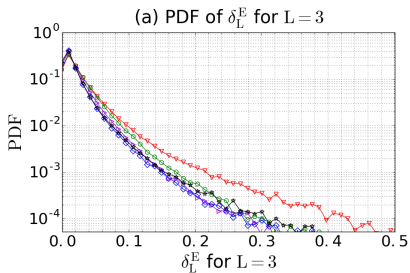
LBM for different fixed $\tau \equiv \tau_0 \rightarrow 0.5$: δ_L^E and $Ma_L = \frac{\langle U_{RMS} \rangle_L}{C_s}$



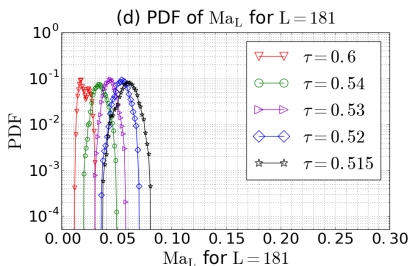
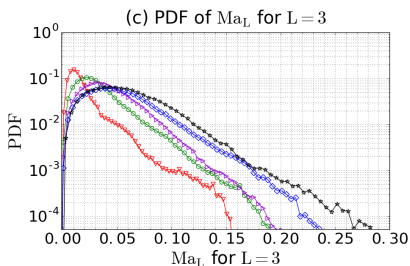
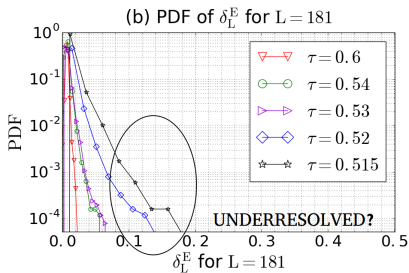
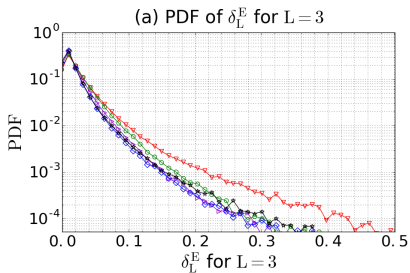
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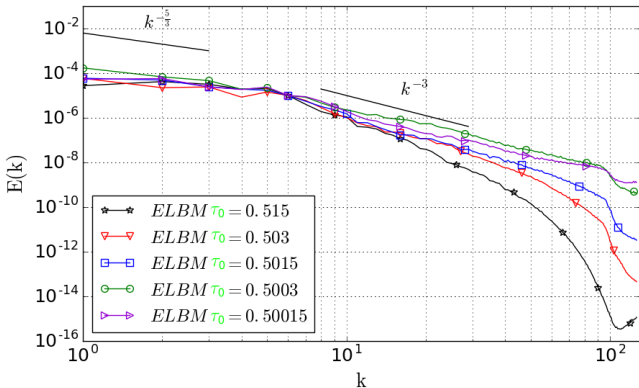
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Entropic LBM for different $\tau_0 \rightarrow 0.5$: Superposed spectrum

$$\tau = \tau^{eff}(\vec{x}, t) = \frac{2\tau_0}{\alpha(f_i(\vec{x}, t))} \text{ \& Constant forcing amplitudes}$$

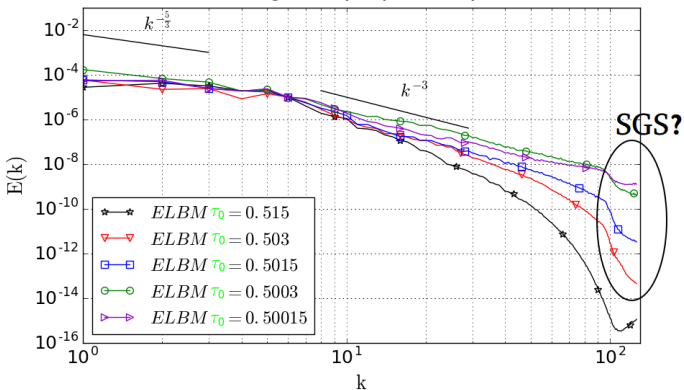
Time-averaged superposed spectrum



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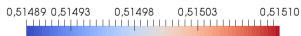
ELBM has the dissipative properties expected from a LES

Entropic LBM for $\tau_0 = 0.515$: Correlation between ω and τ^{eff}

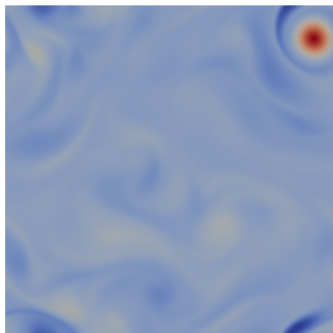
$$\text{Snapshot of } \tau^{eff}(\vec{x}, t) = \frac{2\tau_0}{\alpha(f_i(\vec{x}, t))}$$



Effective relaxation time



$$\text{Snapshot of } \omega = \partial_x u_y - \partial_y u_x$$



Vorticity



Entropic LBM for $\tau_0 = 0.515$: Correlation between ω and τ^{eff}

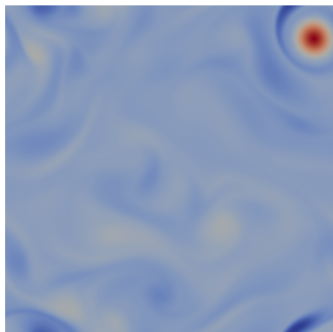
Snapshot of $\tau^{eff}(\vec{x}, t) = \frac{2\tau_0}{\alpha(f_i(\vec{x}, t))}$



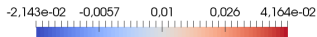
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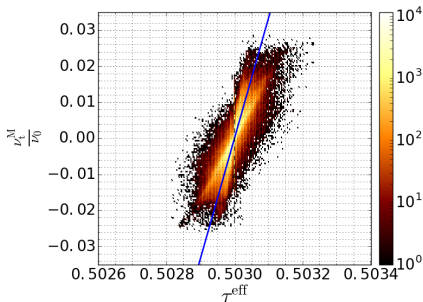


τ^{eff} adapts itself to vorticity peaks

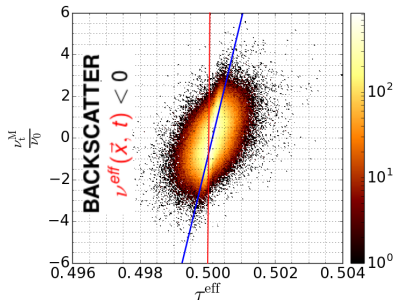
Entropic LBM: Validation of Malaspinas ν_t^M

Is Malaspinas' $\nu_{eff}^M(\vec{x}, t) = \nu_0 + \nu_t^M(\vec{x}, t)$ with $\nu_t^M(\vec{x}, t) \propto -\frac{S_{\theta\kappa} S_{\kappa\gamma} S_{\gamma\theta}}{S_{\lambda\mu} S_{\lambda\mu}}$ an accurate approximation of $\nu_{eff}(\vec{x}, t) = c_s^2 (\tau_{eff}(\vec{x}, t) - 0.5) \Delta t$?

Joint PDF between $\frac{\nu_t^M(\vec{x}, t)}{\nu_0}$ and $\tau_{eff}(\vec{x}, t)$ with $\frac{\nu_t(\vec{x}, t)}{\nu_0} = \frac{\nu_{eff}(\vec{x}, t) - \nu_0}{\nu_0}$ in blue



$$\tau_0 = 0.503$$

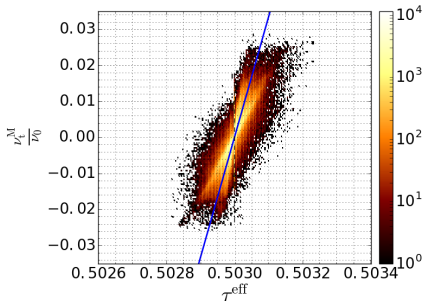


$$\tau_0 = 0.50015$$

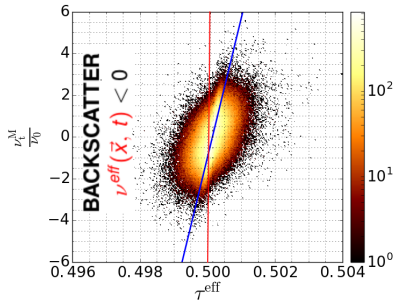
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$$\tau_0 = 0.503$$



$$\tau_0 = 0.50015$$

ν_t^M is indeed a 1st order expansion of ν_t

Conclusions

- ▶ Developed a tool to check numerically the balance of kinetic energy and enstrophy accross scales
- ▶ Applied it to standard LBM:
 - ▶ Hydrodynamics is well recovered at large scales
 - ▶ Recovery at small-scales is less good
 - ▶ Enstrophy balance highlights higher order M_a terms
- ▶ Preliminary results on Entropic LBM:
 - ▶ Dissipative properties as $\tau_0 \rightarrow 0.5$ are as expected for a LES
 - ▶ Malaspinas' $\nu_t^M(\vec{x}, t)$ was numerically shown to be a 1st order expansion of $\nu_t(\vec{x}, t) = \nu^{eff}(\vec{x}, t) - \nu_0$
 - ▶ As $\tau_0 \rightarrow 0.5$, τ^{eff} variance increase and ν^{eff} can become locally negative
- ▶ Systematic analysis of hydrodynamics recovery for Entropic LBM by adding Malaspinas SGS term to the balance equations is on-going

Thank you for your attention!

Any questions?



This project has received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement No' 642069 and was conducted within the activity of ERC Grant No' 339032