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Slide of the Seminar

Entropy Production in Turbulence Paramterisations

Dr. Denny Gohlke

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(P.I. Prof. Luca Biferale)***

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TRR 181 - M4

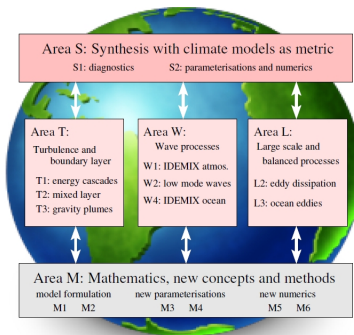
Entropy Production in Turbulence Parameterisations

by

Denny Gohlke

TRR 181 — Energy transfers in Atmosphere and Ocean

www.trr-energytransfers.de



- Cooperation: Universities (Hamburg, Bremen) + German research institutes for meteorology and atmospheric physics
- Goal: understanding of interaction of different dynamical regimes; energetically consistent mathematical models

M4 — Entropy Production in Turbulence Parameterisations

Liebniz-Institute of Atmospheric Physics (Kühlungsborn):
Almut Gassmann and Bastian Sommerfeld (PhD student)

University of Hamburg:
Richard Blender, Denny Gohlke (PhD student) and Valerio Lucarini
(University of Reading)

Main Idea:

Physically consistent representation of
dissipation and backscatter
by use of the **class of Fluctuation Theorems**

Keywords: Non-equilibrium (**entropy production**, 2nd law)
Shell models, turbulence models
Inertial range
Subgrid models

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Class of Fluctuation Theorems — Nonequilibrium Systems

Transient Fluctuation Theorem
(Evans & Searles; 1994, 2002):

$$\frac{p(\bar{\Omega}_t)}{p(-\bar{\Omega}_t)} = e^{\bar{\Omega}_t t}$$

Steady State Fluctuation Theorem
(Gallavotti & Cohen; 1995):

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Crooks Fluctuation Theorem (1999):

$$\frac{p_{A \rightarrow B}(W)}{p_{B \rightarrow A}(-W)} = e^{\beta(W - \Delta F)}$$

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(Seifert 2008):

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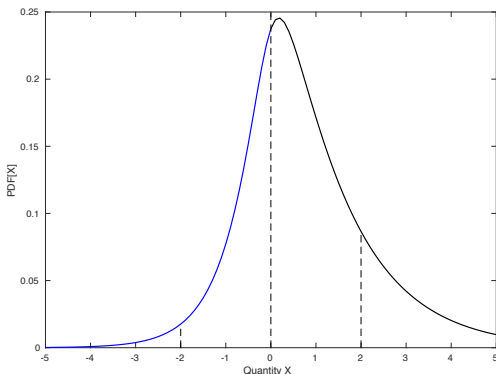
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$$\frac{p(X)}{p(-X)} = e^X, \quad \langle X \rangle \geq 0$$

Gallavotti - Cohen Steady State Fluctuation Theorem

[Gallavotti G, Cohen EGD. 1995. J. Stat. Phys. 80:93170]

Contraction rate in phase space Λ

$$\frac{d}{dt}x = F(x), \Lambda = \nabla \cdot F$$

thermostated:

$$\Lambda \sim EP = Flux/T$$

Chaotic hypothesis: $\bar{\Lambda}_t = \frac{1}{t} \int_0^t \Lambda(\tau) d\tau$ (SRB measure)

Large deviation: $p(\bar{\Lambda}_t) \sim e^{tI(\bar{\Lambda}_t)}$ ($t \rightarrow \infty$)

Time reversibility: $I(\bar{\Lambda}_t) - I(-\bar{\Lambda}_t) = \bar{\Lambda}_t$

$$\lim_{t \rightarrow \infty} \frac{1}{t} \ln \left(\frac{p(\bar{\Lambda}_t)}{p(-\bar{\Lambda}_t)} \right) = \bar{\Lambda}_t$$

Universal slope!

Gallavotti - Cohen Steady State Fluctuation Theorem

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$$\frac{d}{dt}x = F(x), \Lambda = \nabla \cdot F$$

thermostated:

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dissipative:

$$\Lambda < 0$$

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$$\lim_{t \rightarrow \infty} \frac{1}{t} \ln \left(\frac{p(X)}{p(-X)} \right) = X?$$

Quantity?

Universal slope?

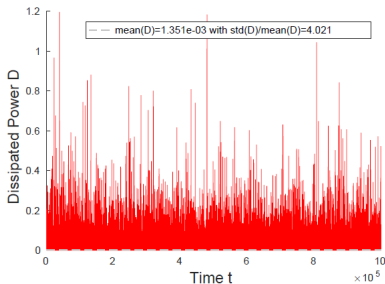
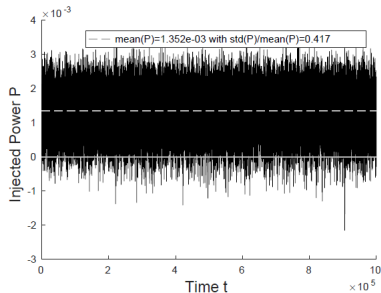
Fluctuation Theorem in dissipative GOY Shell Model?

Gallavotti-Cohen-like

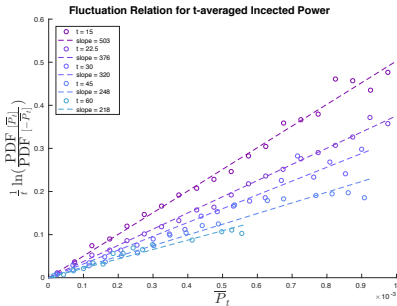
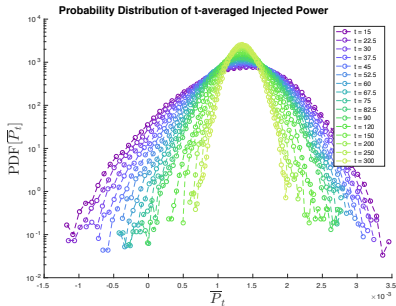
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[Aumaitre S, Fauve S, McNamara S, Poggi P. 2001. Eur. Phys. J. B 19: 449-60]
[Gilbert T. 2004. Europhys. Lett. 67 (2): 172-178]

$$\dot{E}(t) = P(t) + D(t), \quad \langle P \rangle = \langle D \rangle$$

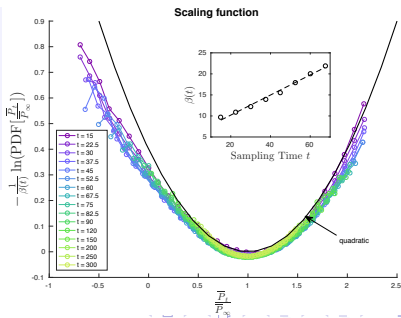


GOY Model — Statistics of t-averaged Injected Power



Distribution of normalised power $Y_t = \frac{\bar{P}_t}{\bar{P}_\infty}$ described by $p(Y_t) \propto e^{-\beta(t)I(Y)}$ with scaling function $I(Y)$ and $\beta = At + B$, leading to **non-asymptotic** relation:

$$\ln\left(\frac{p(Y_t)}{p(-Y_t)}\right) = (At + B)Y_t$$



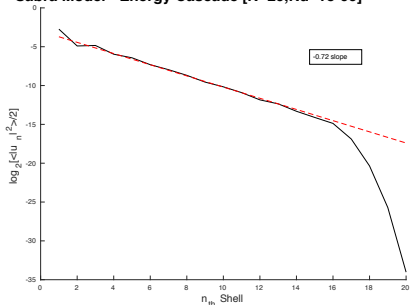
Sabra Model — Statistics of Energy Fluxes Through Shells

$$\frac{d}{dt} \sum_{n=1}^M \frac{1}{2} |u_n|^2 = \underbrace{-k_N \mathfrak{S}(u_M^* u_{M+1}^* u_{M+2}) + \frac{1-\alpha}{2} u_{M-1}^* u_M^* u_{M+1}}_{\Pi_M} + \underbrace{\Re(f u_1^*)}_P - \underbrace{\nu \sum_{n=1}^M k_n^2 |u_n|^2}_D$$

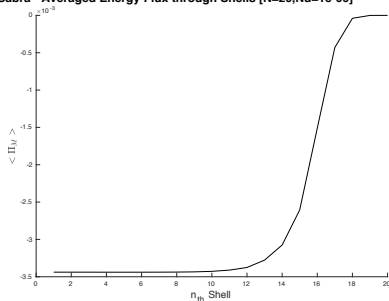
$$\alpha = 1/2 \text{ (3D-like)}$$

$$\langle \Pi_M \rangle = -D \text{ (inertial range)}$$

Sabra Model - Energy Cascade [N=20, Nu=1e-06]



Sabra - Averaged Energy Flux through Shells [N=20, Nu=1e-06]

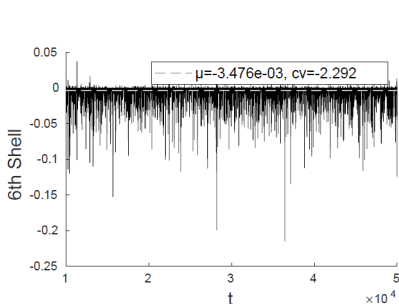


Sabra Model — Statistics of Energy Fluxes Through Shells

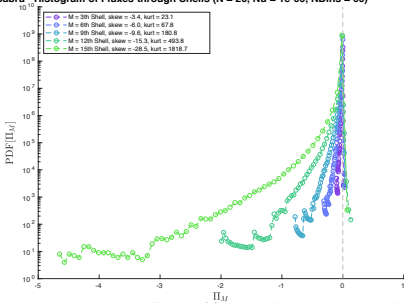
$$\frac{d}{dt} \sum_{n=1}^M \frac{1}{2} |u_n|^2 = \underbrace{-k_N \Im(u_M^* u_{M+1}^* u_{M+2}) + \frac{1-\alpha}{2} u_{M-1}^* u_M^* u_{M+1}}_{\Pi_M} + \underbrace{\Re(f u_1^*)}_P - \underbrace{\nu \sum_{n=1}^M k_n^2 |u_n|^2}_D$$

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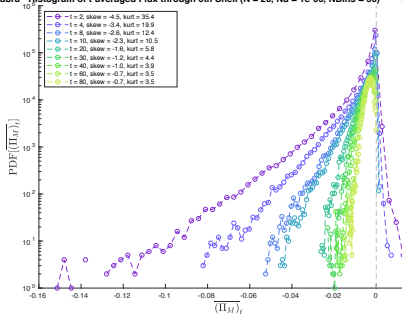


Sabra - Histogram of Fluxes through Shells (N = 20, Nu = 1e-06, NBins = 60)

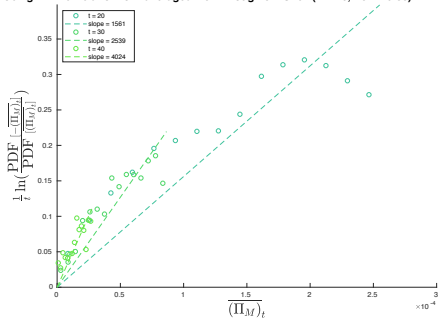


Sabra Model — Statistics of Energy Fluxes Through Shells

Sabra - Histogram of t-averaged Flux through 6th Shell (N = 20, Nu = 1e-06, NBins = 50)



Probing FT-like Relation for t-averaged Flux through 6th Shell (N = 20, Nu = 1e-06)

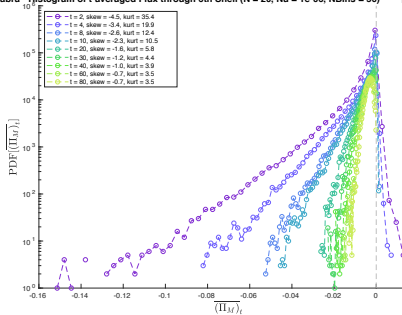


Problems:

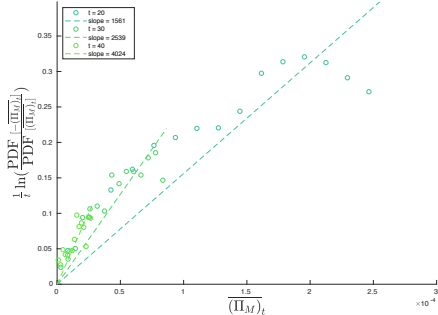
- Detection of large deviations
- Definition of entropy production (incl. temperature)

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Probing FT-like Relation for t-averaged Flux through 6th Shell (N = 20, Nu = 1e-06)



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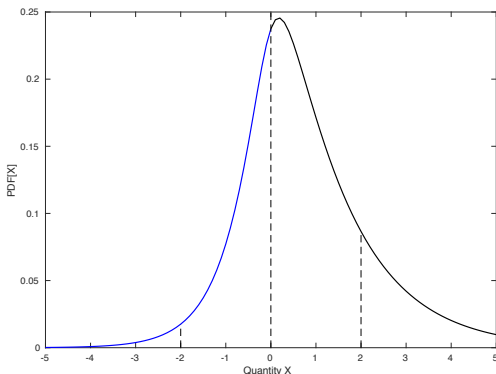
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[Seifert U. 2008. Eur. Phys. J. B 64]

Coupled Langevin System

$$\frac{d}{dt} \vec{u} = \vec{F}(\vec{u}) + \vec{\xi}(t), \quad \langle \vec{\xi}(t) \otimes \vec{\xi}(t') \rangle = 2 \underline{\underline{D}} \delta(t - t')$$

Entropy Change along a Stochastic Trajectory

$$\Delta S_{tot} = \Delta S_{env} + \Delta S_p$$

Concepts from Thermodynamics:

1. From Energy Balance: $\Delta S_{env} = \int_0^\tau dt \frac{d}{dt} \vec{u}(t) \underline{\underline{D}}^{-1} \vec{F}(\vec{u}(t))$
2. Entropy of Particle: $\Delta S_p = -\ln \underbrace{p(\vec{u}(t), t)}_{\text{Fokker-Planck}} \Big|_0^\tau$

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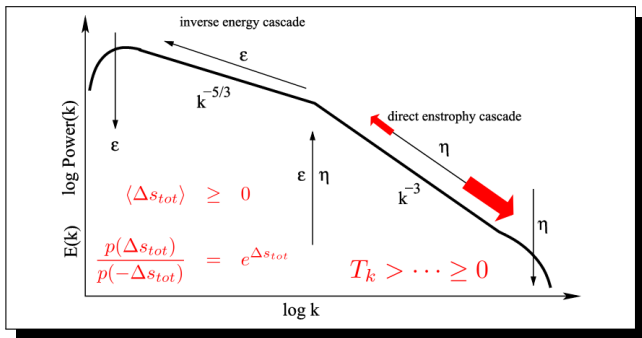
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Enstrophy Cascade and Fluctuations in 2D Turbulence



Vorticity Equation

[Baiesi M, Maes C. 2005. Phys. Rev. E 72]

$$d\omega_k = [F_k(\omega) - \nu k^2 \omega_k] dt + \sqrt{2\gamma_k} dW_k(t)$$

Heat Conduction Problem

$$T_k = \gamma_k / \nu k^2$$

$$\Delta S_{tot} = \sum_k J_k / T_k + \ln p(\omega(0)) / p(\omega(\tau))$$

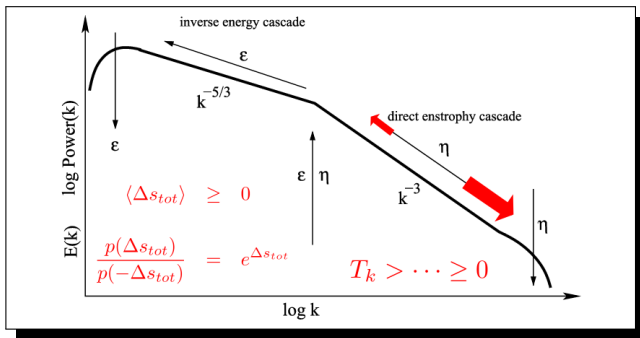
Sabra Shell Model ($k \rightarrow k_n$):

$$\begin{aligned} du_n = & \underbrace{[F_n(u) - \nu k_n^2 u_n]}_{=} dt + \sqrt{2\gamma_n} dW_n(t) \\ & ik_n (u_{n+1}^* u_{n+2} - \frac{\alpha}{2} u_{n-1}^* u_{n+1} + \frac{1-\alpha}{4} u_{n-1} u_{n-2}) \end{aligned}$$

What about ΔS_n ?

[Benzi R, Biferale L, Sbragaglia M. 2004. J. Stat. Phys. 114: 137-154]

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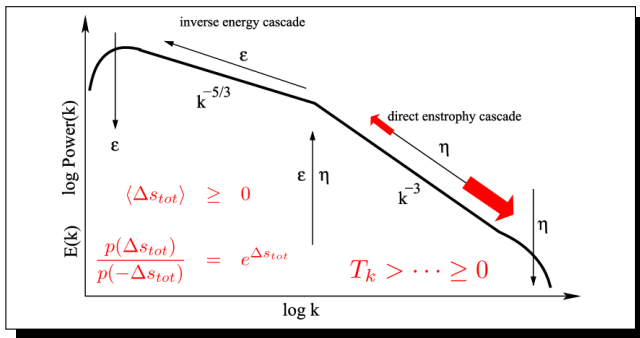
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Reduced Sabra Model — Wouters-Lucarini Parameterisation

$$\frac{d}{dt} u_n = \underbrace{ik_n (u_{n+1}^* u_{n+2} - \frac{\alpha}{2} u_{n-1}^* u_{n+1} + \frac{1-\alpha}{4} u_{n-1} u_{n-2})}_{\text{energy conserving}} + \underbrace{f_n}_{\text{large scales}} - \underbrace{\nu k_n^2 u_n}_{\text{small scales}}$$

[Wouters J. Lucarini V. J. Stat. Mech. (2012): P03003]

$$\text{resolved: } \frac{d}{dt} X = F_X(X) + \epsilon \Psi_X(X, Y)$$

$$\text{unresolved: } \frac{d}{dt} Y = F_Y(Y) + \epsilon \Psi_Y(X, Y)$$

Weak coupling + Chaotic dynamics

$$\rho_{XY, \epsilon}(A(X)) = \rho_{X_{\text{Reduced}}, \epsilon}(A(X)) + \mathcal{O}(\epsilon^3) \quad (\text{SRB measures})$$

$$\text{Reduced: } \frac{d}{dt} X = F_X(X) + \underbrace{\epsilon D(X)}_{=\langle \Psi_X \rangle_{\rho_{Y,0}}} + \underbrace{\epsilon^2 S(X, t)}_{\text{corr}(\Psi_X'; \rho_{Y,0})} + \underbrace{\epsilon^2 M(X, t)}_{=\int_0^\infty h(X(t-s), s) ds}$$

$$h(X, s) = \langle \Psi_Y(X, Y)^T \cdot \partial_Y \Psi_X(f_X^s(X), f_Y^s(Y)) \rangle_{\rho_{Y,0}}$$

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$$h(X, s) = \langle \Psi_Y(X, Y)^T \cdot \partial_Y \Psi_X(f_X^s(X), f_Y^s(Y)) \rangle_{\rho_{Y,0}}$$

Reduced Sabra Model — Wouters-Lucarini Parameterisation

$$\begin{aligned} M(X, t) = & \int_0^\infty \sum_{k,p,q} C_{kpq}^{(1)} f_X^s(X_k(t-s)) X_p(t-s) X_q(t-s) ds \\ & + \int_0^\infty \sum_{p,q} C_{pq}^{(2)} X_p(t-s) X_q(t-s) ds \\ & + \int_0^\infty \sum_{k,r} C_{kr}^{(3)} f_X^s(X_k(t-s)) X_r(t-s) ds \\ & + \int_0^\infty \sum_r C_r^{(4)} X_r(t-s) ds \end{aligned}$$

Weighting terms determined by:

$$\langle \partial_Y \otimes Y^s \rangle_{\rho_{Y,0}}$$

$$\langle Y^s \otimes \partial_Y \otimes Y^s \rangle_{\rho_{Y,0}}$$

$$\langle Y \otimes Y^s \otimes \partial_Y \otimes Y^s \rangle_{\rho_{Y,0}}$$

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