An overview of Fourier pseudospectral methods

Numerical simulation of passive scalar turbulence

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A database of turbulent fields for olfactory search



Incompressible Navier-Stokes turbulence

$$\nabla_i u_i = 0$$

$$\partial_t u_i + \nabla_j (u_i u_j) = -\nabla_i P + \nu \nabla^2 u_i + f_i$$

Velocity **u**(**x**,t) Pressure by density P(**x**,t) External forcing **f**(**x**,t)



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Fourier representation

Employ FFT algorithm

•

$$\hat{\boldsymbol{u}}(\boldsymbol{k},t) = \frac{1}{L^3} \int \mathrm{d}\boldsymbol{x} \boldsymbol{u}(\boldsymbol{x},t) \exp(-\mathrm{i}\boldsymbol{k} \cdot \boldsymbol{x})$$

- Efficient, versatile, easy to parallelize
- Discretize: Wavevector $\mathbf{k} = \mathbf{k} \Delta k$ and position vector $\mathbf{x} = \mathbf{j} \Delta x$

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$$\begin{split} \hat{u}_i(\boldsymbol{k},t) &= \frac{1}{N^d} \sum_{\boldsymbol{j}} u_i(\boldsymbol{j}\Delta x,t) \exp(-\mathrm{i}\boldsymbol{j} \cdot \boldsymbol{k}\Delta x\Delta k) \\ \Delta x \Delta k &= \frac{2\pi}{N} \quad \begin{array}{l} \text{Choice to make FFT work} \\ \text{Domain length } \boldsymbol{L} \text{ discretized using } \boldsymbol{N} \text{ points} \end{split}$$

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• Discrete FFT of Navier-Stokes:

$$k_i \hat{u}_i = 0$$

$$\partial_t \hat{u}_i(\boldsymbol{k}, t) = Q_{ijl}(\boldsymbol{k}) \sum_{\boldsymbol{p}+\boldsymbol{q}=\boldsymbol{k}} \hat{u}_j(\boldsymbol{p}, t) \hat{u}_l(\boldsymbol{q}, t) - \nu \|\boldsymbol{k}\|^2 \hat{u}_i(\boldsymbol{k}, t) + \hat{f}_i(\boldsymbol{k}, t)$$

Convolution sums and related issues

$$\begin{aligned} k_i \hat{u}_i &= 0\\ \partial_t \hat{u}_i(\boldsymbol{k}, t) &= Q_{ijl}(\boldsymbol{k}) \sum_{\boldsymbol{p} + \boldsymbol{q} = \boldsymbol{k}} \hat{u}_j(\boldsymbol{p}, t) \hat{u}_l(\boldsymbol{q}, t) - \nu k^2 \hat{u}_i(\boldsymbol{k}, t) + \hat{f}_i(\boldsymbol{k}, t)\\ Q_{ijl} &= -\mathrm{i}k_l \left(\delta_{ij} - \frac{k_i k_j}{\|\boldsymbol{k}\|^2} \right) \end{aligned}$$

• Lack of analytic solutions: Burgers eq. for particular Q_{ijl} (cf. Lorenz system)

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- Lack of analytic solutions: Burgers eq. for particular Q_{ijl} (cf. Lorenz system)
- Energy and enstrophy transfer (dimensionality-dependent)
- Skewness and intermittency of the velocity gradients
- Products in physical space rather than direct convolution
- Aliasing removal

Building a Fourier pseudospectral code



Aliasing removal by padding

$$\sum_{\boldsymbol{p}+\boldsymbol{q}=\boldsymbol{k}} \hat{u}_j(\boldsymbol{p},t) \hat{u}_l(\boldsymbol{q},t) = \left[u_j(\boldsymbol{x},t) u_l(\boldsymbol{x},t) \right] (\boldsymbol{k},t)$$
$$-N/2 \le k_i < N/2$$

 $-M/2 \le p_i, q_i < M/2$



Aliasing removal by padding



Aliasing removal by padding



- Removal by truncation [2/3 rule] (padding preferable [3/2 rule])
- Alternatives: phase shift, implicit, ..
- ..or combined with some approximations (Rogallo 1977)

Using incompressibility: 2D

$$\begin{split} u_1 &= -\nabla_2 \psi \\ u_2 &= \nabla_1 \psi \\ \partial_t \nabla^2 \psi &= -\nabla_1 \nabla^2 \psi \nabla_2 \psi + \nabla_2 \nabla^2 \psi \nabla_1 \psi + \nu \nabla^2 \nabla^2 \psi \end{split}$$

• All info in one scalar function $\boldsymbol{u} = \boldsymbol{Rot}(\boldsymbol{\psi}\boldsymbol{e}_z)$





- Craya decomposition: *u* actually 2D
- Only traceless part of convective *C_{ij}* matters
- *C_{ij}* traceless: We can spare one FFT (Rogallo 1981, Basdevant 1983)
- 3+6-1 = 8 FFT for one time step



• Unstable (avoid $u^*grad(...)$): $\nabla \cdot (uu) = u \cdot \nabla u$

• Stable, great for helical flows:
$$\nabla \cdot (uu) = \omega \times u + \nabla \frac{u \cdot u}{2}$$

Common types of forcing

• Stochastic: dissipation rate through FDN (Furutsu 1963)

$$\left\langle \hat{f}_{i}(\boldsymbol{k},t)\hat{f}_{j}(\boldsymbol{k}',t')\right\rangle = \delta_{ij}\delta(\boldsymbol{k}+\boldsymbol{k}')\hat{R}(\boldsymbol{k})$$

$$\left\langle u_{i}f_{i}\right\rangle = \nu\left\langle \nabla_{j}u_{i}\nabla_{j}u_{i}\right\rangle = \frac{1}{2}R_{ii}$$

• Keep constant energy (adjust amplitudes within K_F)

$$\hat{\boldsymbol{u}}_F = \hat{\boldsymbol{u}}'_F \sqrt{1 + \frac{E' - \overline{E}}{E'_F}}$$

• Statistically steady input energy input rate

$$\hat{\boldsymbol{f}}(\boldsymbol{k},t) = \epsilon \frac{\hat{\boldsymbol{u}}(\boldsymbol{k},t)}{\sum_{K_F} \|\hat{\boldsymbol{u}}\|^2}, \, \boldsymbol{k} \in K_F$$

Time integration: Exponential Runge Kutta

• Integrating factor

 $\partial_t \left(\boldsymbol{u} \exp(\nu k^2 t) \right) = \exp(\nu k^2 t) \boldsymbol{F}[\boldsymbol{u}, t]$

• Integrate and discretize rhs using Lagrangian polynomials

Time integration: Exponential Runge Kutta

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- Integrate and discretize rhs using Lagrangian polynomials
- E.g., second-order exponential Runge Kutta (Hochbruck 2010)

$$\begin{aligned} \boldsymbol{u}_1' &= \varphi_0 \boldsymbol{u}_0 + \varphi_1 \boldsymbol{F}_0 \\ \boldsymbol{u}_1 &= \varphi_0 \boldsymbol{u}_0 + (\varphi_1 - \varphi_2) \boldsymbol{F}[\boldsymbol{u}_0, t] + \varphi_2 \boldsymbol{F}[\boldsymbol{u}_1', t + \Delta t] \\ z &= -\nu k^2 \Delta t \quad \varphi_0 = e^z \quad \varphi_1 = \frac{e^z - 1}{z} \quad \varphi_2 = \frac{e^z - z - 1}{z^2} \end{aligned}$$

FFT pencil parallelization

- P3DFFT (Pekurovsky 2012)
- Transpositions: MPI Alltoall
- Local FFTW or, recently, cuFFT (now available 3D parallel)



Parallel I/O

- Collective read/write with derived data-types (domain slices)
- avoid write/read_at, non-shared headers (will lock the file)
- MPI_File_write_all(...)



Particle phase: Non-Uniform FFT

• Connect Eulerian fields and point-like particles

$$\boldsymbol{C}(\boldsymbol{x},t) = \sum_{p} c_{p}(t)\delta(\boldsymbol{x}-\boldsymbol{x}_{p})$$

- Trick: take convolution in physical space..
- ..and then remove it in Fourier

$$\hat{\boldsymbol{C}}(\boldsymbol{x},t) \simeq \widehat{\boldsymbol{C}(\boldsymbol{x},t) * B(\boldsymbol{x})}$$
 $\hat{\hat{B}}(\boldsymbol{k})$

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• MPI Neighbor collectives for particle parallelization

DNS code in a nutshell

- Built-in P3DFFT
- Basdevant trick
- MPI parallel IO
- No external libraries
- NUFFT for field interpolation
- NUFFT for particles feedback
- MPI Neighbor collectives
- Weakly compressible module for particle combustion



Extension to weakly compressible



- Immersed boundaries, particles, and much more
- Iron particle combustion
- Generates large heat and mass fluxes
- Multiscale interaction with turbulence

2D passive scalar turbulence model

Hyperviscosity
degree 8

$$\partial_t u_i + \nabla_j (u_j u_i) = -\nabla_i P + \nu(\zeta) \nabla^{2\zeta} u_i - u_i / \mathcal{T}_u + f_i$$

Large-scale Gaussian
damping forcing
 $\partial_t \theta + \nabla_j (u_j \theta) = \kappa \nabla^2 \theta - \theta / \mathcal{T} + RG (x - x_s)$
Large-scale Smoothed

Large-scale Smoothed damping point-like source

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Large-scale Gaussian
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 $\partial_t \theta + \nabla_j (u_j \theta) = \kappa \nabla^2 \theta - \theta / \mathcal{T} + RG (\boldsymbol{x} - \boldsymbol{x}_s)$
Large-scale Smoothed
damping point-like
source
 $G(\boldsymbol{x}) = \frac{1}{2\pi\sigma^2} \exp\left(-\|\boldsymbol{x}\|^2 / (2\sigma^2)\right)$

$$\langle f_i(\boldsymbol{x},t) f_j(\boldsymbol{x}',t') \rangle = \delta_{ij} \delta(t-t') C(\boldsymbol{x}-\boldsymbol{x}')$$

Gaussian forcing at small scales: **inverse cascade** Gaussian-like correlation function *C*







2D passive scalar: average concentration



• Mask periodicity effects on ~10% domain size

 $\theta(\boldsymbol{x},t) = \prod_{i} \theta'(\boldsymbol{x},t) \tanh(S\boldsymbol{e}_{i} \cdot (\boldsymbol{x} - \boldsymbol{x}_{b,i}))$

- Average scalar concentration: ~ 5 to <0.1 (decrease damping time)
- Resolved time series at the red points



2D passive scalar: single-time statistics



2D passive scalar: correlations



• Hyperviscosity and mask effects

Scalar time series zero mean wind





Scalar time series with mean wind





Scalar time series double mean wind





What's next

